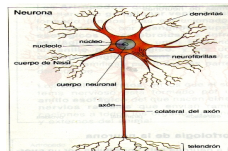
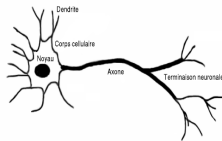


PDEs for neural networks analysis and behaviour

Benoît Perthame



- I. The single neuron,

- II. Networks (Wilson-Cowan, Integrate & Fire networks)

- III. Networks, time elapsed models
 - III. 1. Desynchronisation

 - III. 2. Spontaneous activity

 - III. 3. Relation to I&F



Based on K. Pakdaman, J. Champagnat, J.-F. Vibert

- s represents the time elapsed since the last discharge
- $n(s, t)$ probability of finding a neuron in 'state' s at time t
- $N(t)$ = activity of the network

$$\frac{\partial n(s, t)}{\partial t} + \overbrace{\frac{\partial n(s, t)}{\partial s}}^{\text{elapsed time advances}} + \overbrace{r(s, bN(t)) n(s, t)}^{\text{firing neurons}} = 0,$$

$$N(t) := n(s = 0, t) = \underbrace{\int_0^{+\infty} r(s, bN(t)) n(s, t) ds}_{\text{reset of firing neurons}}$$

$$\begin{cases} \frac{\partial n(s,t)}{\partial t} + \frac{\partial n(s,t)}{\partial s} + r(s) n(s,t) = 0, \\ N(t) := n(s=0, t) = \int_0^{+\infty} b(s) n(s,t) ds \end{cases}$$

We want a steady state (compute the growth rate)

Long standing equation (Feller, renewal process)

$$\begin{cases} \frac{\partial n(s,t)}{\partial t} + \frac{\partial n(s,t)}{\partial s} + [\lambda + r(s)] n(s,t) = 0, \\ N(t) := n(s=0, t) = \int_0^{+\infty} b(s) n(s,t) ds \end{cases}$$

We want a steady state

$$\bar{n}(s) = \bar{n}(0)e^{-R(s)-\lambda s}, \quad R(s) = \int_0^s r$$

In the boundary condition

$$\bar{n}(0) = \bar{n}(0) \int_0^{\infty} b(s)e^{-R(s)-\lambda s} ds$$

Find λ such that

$$\int_0^{\infty} b(s)e^{-R(s)-\lambda s} ds = 1$$

$$\begin{cases} \frac{\partial n(s,t)}{\partial t} + \frac{\partial n(s,t)}{\partial s} + [\lambda + r(s)] n(s,t) = 0, \\ N(t) := n(s=0, t) = \int_0^{+\infty} b(s) n(s,t) ds \end{cases}$$

$$\bar{n}(s) = \bar{n}(0) e^{-R(s)-\lambda s}, \quad R(s) = \int_0^s r$$

Find λ such that

$$\int_0^{\infty} b(s) e^{-R(s)-\lambda s} ds = 1$$

When $b(\cdot) = r(\cdot)$ then $\lambda = 0$

$$\begin{cases} \frac{\partial n(s,t)}{\partial t} + \frac{\partial n(s,t)}{\partial s} + [\lambda + r(s)] n(s,t) = 0, \\ N(t) := n(s=0, t) = \int_0^{+\infty} b(s) n(s,t) ds \end{cases}$$

$$\bar{n}(s) = \bar{n}(0)e^{-R(s)-\lambda s}, \quad \int_0^{\infty} b(s)e^{-R(s)-\lambda s} ds = 1$$

Prove that

$$n(s, t) \rightarrow \bar{n}(s)$$

Use Generalized Relative Entropy

$$\begin{cases} \frac{\partial n(s,t)}{\partial t} + \frac{\partial n(s,t)}{\partial s} + [\lambda + r(s)] n(s,t) = 0, \\ N(t) := n(s=0, t) = \int_0^{+\infty} b(s) n(s,t) ds \\ \frac{\partial \phi(s)}{\partial s} - [\lambda + r(s)] \phi(s) = -\phi(0)b(s) \end{cases}$$

We have

$$\frac{d}{dt} \int n(s,t) \phi(s) = 0$$

Usually for a conservative equation $\phi \equiv 1$

For example, when $b(\cdot) = r(\cdot)$, $\lambda = 0$.

Then the relative entropy is $\int_0^{\infty} \bar{n}(s) H\left(\frac{n(s,t)}{\bar{n}(s)}\right) ds$

$$\begin{cases} \frac{\partial n(s,t)}{\partial t} + \frac{\partial n(s,t)}{\partial s} + [\lambda + r(s)] n(s,t) = 0, \\ N(t) := n(s=0, t) = \int_0^{+\infty} b(s) n(s,t) ds \end{cases}$$

$$\frac{\partial \phi(s)}{\partial s} - [\lambda + r(s)] \phi(s) = -\phi(0)b(s)$$

The Generalized Relative Entropy is

$$E(t) := \int_0^{\infty} \phi(s) \bar{n}(s) H\left(\frac{n(s,t)}{\bar{n}(s)}\right) ds$$

$$\begin{cases} \frac{\partial n(s,t)}{\partial t} + \frac{\partial n(s,t)}{\partial s} + [\lambda + r(s)] n(s,t) = 0, \\ N(t) := n(s=0, t) = \int_0^{+\infty} b(s) n(s,t) ds \\ \frac{\partial \phi(s)}{\partial s} - [\lambda + r(s)] \phi(s) = -\phi(0)b(s) \end{cases}$$

The Generalized Relative Entropy is

$$E(t) := \int_0^{\infty} \phi(s) \bar{n}(s) H\left(\frac{n(s,t)}{\bar{n}(s)}\right) ds$$

$$\frac{dE(t)}{dt} = -D_H(t)\phi(0)$$

$$D_H(t) = \int b(s) \bar{n}(s) H\left(\frac{n(s,t)}{\bar{n}(s)}\right) ds - H\left(\int b(s) \bar{n}(s) \frac{n(s,t)}{\bar{n}(s)} ds\right)$$

> 0 by Jensen inequality

And a Poincaré inequality?

$$\int_0^\infty \phi(s) \bar{n}(s) |u(s)| ds \leq \nu \left[\int_0^\infty b(s) \bar{n}(s) |u(s)| ds - \left| \int_0^\infty b(s) \bar{n}(s) u(s) ds \right| \right]$$

when

$$\int_0^\infty b(s) \bar{n}(s) ds = 1, \quad \int_0^\infty u(s) \bar{n}(s) \phi(s) ds = 0$$

Theorem

If $b(\cdot) \geq \nu \phi(\cdot)$, then the Poincaré inequality holds true.

Proof

$$\begin{aligned} \left| \int_0^\infty b(s) \bar{n}(s) u(s) ds \right| &= \left| \int_0^\infty [b(s) - \nu \phi(s)] \bar{n}(s) u(s) ds \right| \\ &\geq \int_0^\infty [b(s) - \nu \phi(s)] \bar{n}(s) |u(s)| ds \end{aligned}$$

And a Poincaré inequality?

$$\nu \int_0^\infty \phi(s) \bar{n}(s) |u(s)| ds \leq \int b(s) \bar{n}(s) |u(s)| ds - \left| \int_0^\infty b(s) \bar{n}(s) u(s) ds \right|$$

when

$$\begin{aligned} \int b(s) \bar{n}(s) &= 1, & \int u(s) \bar{n}(s) \phi(s) ds &= 0 \\ \left| \int_0^\infty b(s) \bar{n}(s) u(s) ds \right| &= \left| \int_0^\infty [b(s) - \nu \phi(s)] \bar{n}(s) u(s) ds \right| \\ &\geq \int_0^\infty [b(s) - \nu \phi(s)] \bar{n}(s) |u(s)| ds \\ &= \int_0^\infty b(s) \bar{n}(s) |u(s)| ds - \int_0^\infty \nu \phi(s) \bar{n}(s) |u(s)| ds \end{aligned}$$

The GRE is robust

$$\begin{cases} \frac{\partial n(s,t)}{\partial t} + \frac{\partial n(s,t)}{\partial s} + r(s,t) n(s,t) = 0, \\ N(t) := n(s=0, t) = \int_0^{+\infty} b(s,t) n(s,t) ds \end{cases}$$

Assume that $r(s, t)$, $b(s, t)$ are time periodic

$$\begin{cases} \frac{\partial n(s,t)}{\partial t} + \frac{\partial n(s,t)}{\partial s} + [\lambda_F + r(s,t)] n(s,t) = 0, \\ N(t) := n(s=0, t) = \int_0^{+\infty} b(s,t) n(s,t) ds \end{cases}$$

$r(s,t)$, $b(s,t)$ are time periodic

$$\begin{cases} \frac{\partial n(s,t)}{\partial t} + \frac{\partial n(s,t)}{\partial s} + [\lambda_F + r(s,t)] n(s,t) = 0, \\ N(t) := n(s=0, t) = \int_0^{+\infty} b(s,t) n(s,t) ds \end{cases}$$

$r(s,t)$, $b(s,t)$ are time periodic

$$\begin{cases} \frac{\partial \bar{n}(s,t)}{\partial t} + \frac{\partial \bar{n}(s,t)}{\partial s} + [\lambda_F + r(s,t)] \bar{n}(s,t) = 0, \\ \bar{N}(t) := \bar{n}(s=0, t) = \int_0^{+\infty} b(s,t) \bar{n}(s,t) ds \end{cases}$$

$$\frac{\partial \bar{\phi}(s,t)}{\partial t} + \frac{\partial \bar{\phi}(s,t)}{\partial s} - [\lambda_F + r(s,t)] \bar{\phi}(s,t) = \phi(0,t)b(s,t),$$

with

$$\bar{n}(s,t) > 0, \bar{\phi}(s,t) > 0 \text{ time periodic}$$

$$\begin{cases} \frac{\partial n(s,t)}{\partial t} + \frac{\partial n(s,t)}{\partial s} + [\lambda_F + r(s,t)] n(s,t) = 0, \\ N(t) := n(s=0, t) = \int_0^{+\infty} b(s,t) n(s,t) ds \end{cases}$$

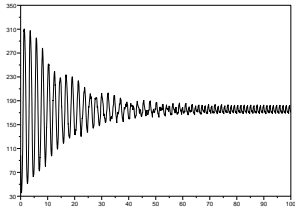
The Generalized Relative Entropy is

$$E(t) := \int_0^{\infty} \bar{\phi}(s,t) \bar{n}(s,t) H\left(\frac{n(s,t)}{\bar{n}(s,t)}\right) ds$$

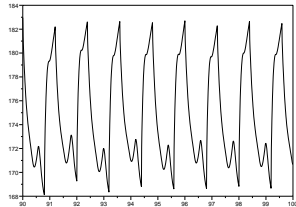
$$\frac{dE(t)}{dt} = -D_H(t) \phi(0, t)$$

$$D_H(t) = \int b(s,t) \bar{n}(s,t) H\left(\frac{n(s,t)}{\bar{n}(s,t)}\right) ds - H\left(\int b(s,t) \bar{n}(s,t) \frac{n(s,t)}{\bar{n}(s,t)} ds\right)$$

> 0 by Jensen inequality



global simulation



zoom on the last periods

Total cell density $\int_0^\infty n(s, t) ds$ as a function of time

$$\begin{aligned}
 & \text{elapsed time advances} \\
 & \frac{\partial n(s,t)}{\partial t} + \frac{\partial n(s,t)}{\partial s} + \overbrace{r(s, bN(t)) n(s,t)}^{\text{firing neurons}} = 0, \\
 N(t) := n(s=0, t) &= \underbrace{\int_0^{+\infty} r(s, bN(t)) n(s,t) ds}_{\text{reset of firing neurons}},
 \end{aligned}$$

Theorem For small or large connectivity ($b > 0$ small or large) then desynchronization still holds

$$n(s, t) \xrightarrow[t \rightarrow \infty]{} P_b(s)$$

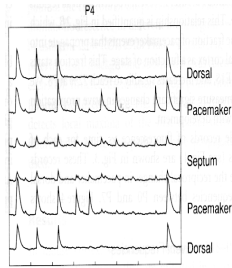
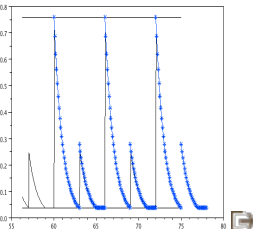
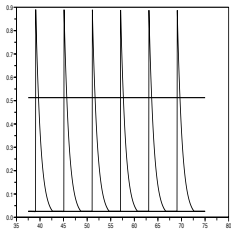
$$\begin{aligned}
 & \text{elapsed time advances} \quad \text{firing neurons} \\
 & \frac{\partial n(s,t)}{\partial t} + \frac{\partial n(s,t)}{\partial s} + r(s, bN(t)) n(s,t) = 0, \\
 & N(t) := n(s=0, t) = \underbrace{\int_0^{+\infty} r(s, bN(t)) n(s,t) ds}_{\text{reset of firing neurons}}
 \end{aligned}$$

Theorem For small or large connectivity ($b > 0$ small or large) then desynchronization still holds

$$n(s, t) \xrightarrow[t \rightarrow \infty]{} P_b(s)$$

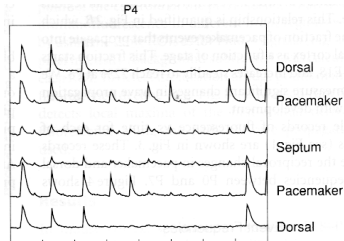
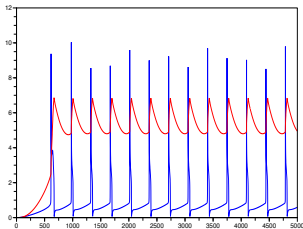
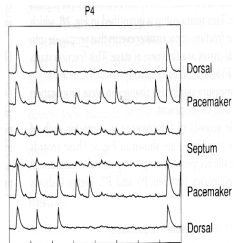
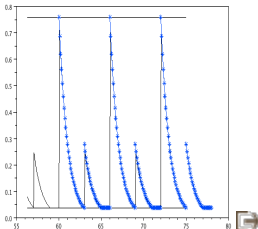
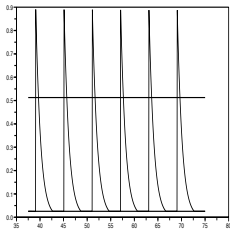
- In the middle range connectivity there are several periodic solutions (analytic forms of solutions),
- These are stable (observed numerically).

$$\begin{aligned}
 & \text{elapsed time advances} \\
 & \frac{\partial n(s,t)}{\partial t} + \frac{\partial n(s,t)}{\partial s} + \overbrace{r(s, bN(t)) n(s,t)}^{\text{firing neurons}} = 0, \\
 & N(t) := n(s=0, t) = \underbrace{\int_0^{+\infty} r(s, bN(t)) n(s,t) ds}_{\text{reset of firing neurons}}
 \end{aligned}$$



Right : Conhaim et al (2011) J. of physiology 589(10) 2529-2541.

Comparison with I&F



Two equations

Time elapsed

$$\begin{cases} \partial_t n + \partial_s n + r(s)n = 0, \\ n(s=0, t) = \int_0^\infty r(s)n(s, t)ds =: N(t), \\ n(s, 0) = n^0(s), \end{cases}$$

LIF

$$\begin{cases} \partial_t \hat{n} + \partial_v [(-v + \sigma)\hat{n}] - a\partial_v^2 \hat{n} = N(t)\delta(v - V_R), \\ \hat{n}(V_F, t) = 0, \quad -a\partial_v \hat{n}(V_F, t) = N(t). \end{cases}$$

From G. Dumont and J. Henry (linear equations)

Theorem There are functions $r(s)$, $q(v, s)$ such that, for solutions of

$$\begin{cases} \partial_t n + \partial_s n + r(s)n = 0, \\ n(s=0, t) = \int_0^\infty r(s)n(s, t)ds =: N(t), \\ n(s, 0) = n^0(s), \end{cases}$$

then

$$\hat{n}(v, t) = \int_0^\infty q(v, s)n(s, t)ds$$

satisfies, with the same $N(t)$,

$$\begin{cases} \partial_t \hat{n} + \partial_v [(-v + \sigma)\hat{n}] - a\partial_v^2 \hat{n} = N(t)\delta(v - V_R), \\ \hat{n}(V_F, t) = 0, \quad -a\partial_v \hat{n}(V_F, t) = N(t). \end{cases}$$

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D. Smets



D. Salort



K. Pakdaman



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